# THE BRENIER-SCHRÖDINGER PROBLEM WITH RESPECT TO FELLER SEMIMARTINGALES AND NON-LOCAL HAMILTON-JACOBI-BELLMAN EQUATIONS 

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#### Abstract

Motivated by a problem from incompressible fluid mechanics of Brenier [Bre89], and its recent entropic relaxation by [ACLZ20], we study a problem of entropic minimization on the path space when the reference measure is a generic Feller semimartingale. We show that, under some regularity condition, our problem connects naturally with a, possibly non-local, version of the Hamilton-JacobiBellman equation. Additionally, we study existence of minimizers when the reference measure in a Ornstein-Uhlenbeck process.


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## 1. INTRODUCTION

1.1. Main result. We study the so-called Brenier-Schrödinger problem with respect to a reference measure $R$, that is the law of a Feller martingale on $[0,1]$. This problem consists in minimizing the relative entropy of the law P of another process $Y$ under the marginal constraints:
(i) $Y_{t} \sim \mu_{t}$, for all $t \in[0,1]$, where $\left(\mu_{t}\right)$ is the data of a family of probability measure.
(ii) $\left(Y_{0}, Y_{1}\right) \sim \pi$, where $\pi$ is the data of a coupling of the endpoints.

By definition of the relative entropy, every minimizer is absolutely continuous with respect to R. It particular, P is also the law of a semimartingale, and the log-density process $Z$ is a semimartingale. Actually, conditionally on $\left\{X_{0}=x\right\}, Z$ is of the form

$$
Z_{t}=A_{t^{-}}+\psi_{t}^{x}\left(X_{t}\right)
$$

where $A$ is an additive functional, and $(t, z) \mapsto \psi_{t}^{x}(z)$ is some function. We say that the minimizer P is regular provided the associated $\psi^{x}$ is $\mathscr{C}^{1}$ in time, $\mathscr{C}^{2}$ in space, and the semimartingale $\psi^{x}(X)$ is special (see definitions below). We recall that a semimartingale is special as soon as it has bounded jumps. In particular, every continuous semimartingale is special. Our main result is as follows.

Theorem. Consider a regular solution P of the aforementioned problem, and the associated $\psi^{x}$. Then, the additive functional $A$ is absolutely continuous. In particular, there exists a function $p:[0,1] \times \mathbb{R}^{n} \rightarrow$ $\mathbb{R}$ such that

$$
\begin{equation*}
A_{t}=\int_{0}^{t} p_{t}\left(X_{s}\right) \mathrm{d} s, \quad t \in[0,1] \tag{1.1}
\end{equation*}
$$

Moreover, $\psi^{x}$ is a solution to the generalized Hamilton-Jacobi-Bellman equation

$$
\begin{equation*}
\partial_{t} \psi^{x}+\mathrm{e}^{-\psi_{t}^{x}} \mathbf{A}\left(\mathrm{e}^{\psi_{t}^{x}}\right)+p_{s}=0 \tag{1.2}
\end{equation*}
$$

where $\mathbf{A}$ denotes the Markov generator of R .
Remark 1.1. Let us make some comments on the result:
(a) When R is a Markov diffusion, then $\mathbf{A} u(x)=\frac{1}{2} c(x) \cdot \nabla^{2} u(x)+b \cdot \nabla u(x)$, for some diffusion matrix $c$ and drift vector $b$. By the chain rule (1.2) becomes

$$
\partial_{t} \psi_{t}^{x}+\mathbf{A} \psi^{x}+\frac{1}{2} \nabla \psi_{t}^{x} \cdot c \nabla \psi_{t}^{x}+p_{t}=0
$$

which is the usual Hamilton-Jacobi-Bellman equation with pressure.
(b) The function $p$ is interpreted as a pressure field.
(c) In the classical Hamilton-Jacobi-Bellman equation, the term $\mathbf{A}$ would have a negative sign. This phenomenon also has been observed in [ACLZ20]. It can be circumvent by considering the potential $\psi$ associated to the time-reversed measure.

Verifying that there exist regular minimizers is an arduous task, that, in the Brownian case amounts to solve the Navier-Stokes equation (see [ACLZ20] for details). We do not take on such an accomplishment. More modestly, additionally to our main theorem, we show that there exists a minimizer, possibly non-regular, when $R$ is the law of an Ornstein-Uhlenbeck process.
1.2. Motivations and connections to the literature. Understanding what happens to the BrenierSchrödinger problem for general semimartingales, possibly with a jump part is the main motivation for this paper. The Brenier-Schrödinger problem, defined in [ACLZ20], is a relaxation of Brenier's approach [Bre89] to incompressible perfect fluids and Euler equations. This generalization, which can be seen as the entropic relaxation of Brenier original problem, aims at modelling viscous fluid dynamics. The achievements of [ACLZ20] are threefold.
(i) When the reference measure R is Markovian, they study the general shape of minimizers.
(ii) Whenever the reference measure R is the law of a reversible Brownian motion on $\mathbb{R}^{n}$ or $(\mathbb{R} \backslash \mathbb{Z})^{n}$, they show that minimizers of a certain form give rise to solution to the Navier-Stokes equation.
(iii) On the torus, they show that minimizers always exists, whenever $\mu_{t}=$ vol for all $t$, and the coupling $\pi$ has finite entropy relatively to the $\mathrm{R}_{01}$.
The two last points are extended to the compact manifold setting in [GH22]. Let us stress that the particular form of the minimizer needed in [ACLZ20] is stronger than our regularity condition. Indeed, they need to assume the existence of the function $p$ such that (1.1) holds, while (1.1) follows from our analysis. Thus, even in the purely continuous case, our result is less restrictive than that of [ACLZ20].

In the Brownian case, when the noise tends to 0, one recovers in the limit Brenier's original problem of minimization of the kinetic energy, as illustrated in [BM20]. In our possibly non-local setting, understanding the appropriate notion of small noise limit, and whether there exists a equivalent of Brenier's problem is an interesting question that could be explored in future works.
1.3. Outline of the proof. Our proof follows the lines of [ACLZ20]. We use, on the one hand, Itō's formula, and on the other hand, Girsanov theorem to give two different representations of the semimartingale $Z$. At the technical level, since $Z$ is a special semimartingale, it thus admits a unique decomposition as a sum of a local martingale, and predictable process of finite variations. This allows us to identify the two different decompositions and conclude. In the continuous case, as for the Brownian setting of [ACLZ20], the semimartingale is always special and this assumption is unnecessary.

## 2. REMINDERS AND NOTATION

2.1. Semimartingales and their characteristics. We refer to [JS03] for more details.
2.1.1. Path space. Let $\Omega$ denote the space of right-continuous with left-limit paths from $[0,1]$ to $\mathbb{R}^{n}$. The canonical process on $\Omega$ is denoted by $X$, that is

$$
X_{t}(\omega):=\omega_{t}, \quad \omega \in \Omega, t \in[0,1]
$$

The associated canonical filtration defined by

$$
\mathfrak{F}_{t}:=\bigcap_{s>t} \sigma\left(X_{r}: 0 \leq r \leq s\right), \quad t \in[0,1]
$$

Conveniently for $\mathrm{P} \in \mathscr{P}(\Omega)$, we write $\mathrm{P}_{t}$ for the law of $X_{t}$ under P , $\mathrm{P}_{t s}$ for the law of $\left(X_{t}, X_{s}\right)$ under P , and so on.
2.1.2. Processes. An adapted process is a mapping $Y: \Omega \rightarrow \Omega$ such that $Y_{t}$ is $\mathfrak{F}_{t}$-measurable for all $t \in[0,1]$. By definition, all our processes are right-continuous with left-limit. For such $Y$, its jump process is

$$
\Delta Y_{t}:=Y_{t}-Y_{t-}, \quad t \in[0,1] .
$$

We say that $Y$ is predictable whenever seen as the mapping

$$
[0,1] \times \Omega \ni(t, \omega) \mapsto Y_{t}(w)
$$

it is measurable with respect to the $\sigma$-algebra generated by $F \times(t, s]$ for $s<t$ and $F \in \mathfrak{F}_{s}$. A process of finite variation is an adapted process $A=B-C$ with $B$ and $C$ two non decreasing adapted processes. In that case the variation of $A$ is the process $B+C$. We also consider generalized adapted (resp. predictable) processes of the form $W: \Omega \times \mathbb{R}^{n} \rightarrow \Omega$.
2.1.3. Martingales \& semimartingales. Fix $\mathrm{P} \in \mathscr{P}(\Omega)$. A martingale (or a P-martingale to emphasize the dependence on $P$ ) is an adapted process such that:

$$
\mathbf{E}_{\mathrm{P}}\left[M_{t} \mid \mathfrak{F}_{s}\right]=M_{s}, \quad 0 \leq s \leq t \leq 1
$$

We say that $X$ is a semimartingale under P provided

$$
\begin{equation*}
X=X_{0}+M+A \tag{2.1}
\end{equation*}
$$

where $M$ is a P -martingale, and $A$ is a finite variation process ( P -almost surely). The decomposition (2.1) is not unique.
2.1.4. Special semimartingales. If $A$ can be chosen predictable in (2.1), we say that the semimartingale is special. By [JS03, pp. I.4.23-24], a semimartingale is special if and only if, in (2.1), $A$ can be chosen with P-integrable variation, and that a martingale is special as soon as its jump are bounded.
2.1.5. Canonical jumps measure. We consider the canonical jumps process and the canonical jumps measure

$$
\begin{aligned}
& \Delta X_{s}:=X_{s}-X_{s-}, \quad s \in[0,1] \\
& \mu_{t}:=\sum_{s \leq t} \delta_{\left(s, \Delta X_{s}\right)} \mathbf{1}_{\Delta X_{s} \neq 0}
\end{aligned}
$$

The measure $\mu$ is a random element of $\mathcal{M}\left([0,1] \times \mathbb{R}^{d}\right)$, the set of Borel measures on $[0,1] \times \mathbb{R}^{d}$. For a, possibly random, measure $\lambda \in \mathcal{M}\left([0,1] \times \mathbb{R}^{n}\right)$, and a generalized process $W$, we write

$$
(W * \lambda)_{t}:=\int_{0}^{t} \int_{\mathbb{R}^{n}} W_{s}(y) \lambda(\mathrm{d} s \mathrm{~d} y), \quad t \in[0,1] .
$$

2.1.6. Compensator of the jumps measure. For all $P \in \mathscr{P}(\Omega)$, there exists a unique predictable measure $\nu$ such that

$$
\left[\mathbf{E}_{\mathrm{P}}[|W| * \mu]<\infty\right] \Rightarrow\left[\mathbf{E}_{\mathrm{P}}[|W| * \nu]<\infty, \text { and } W * \mu-W * \nu \text { is a P-martingale }\right] .
$$

We call $\nu$ the compensator of $\mu$ (under P ).
2.1.7. Characteristics of a semimartingale. We fix $h(y):=y \mathbf{1}_{|y| \leq 1}$. The big-jumps removed version of $X$ is

$$
X^{h}:=X-(y-h) * \mu .
$$

If $X$ is a semimartingale under P , then $X^{h}$ is a special semimartingale. Thus, there exists a unique decomposition

$$
X^{h}=X_{0}+M^{h}+B^{h}
$$

with $M^{h}$ a martingale, and $B^{h}$ predictable. We call $\left(B^{h}, C, v\right)$ the characteristics of the semimartingale, where $C$ is the quadratic variation of the continuous part of $X$, and $\nu$ the compensator.

Remark 2.1. In general, it is not known whether a semimartingale with prescribed characteristics exists, nor if the characteristics characterise $P$.
2.1.8. Representation of semimartingales with given characteristics. Let R be a semimartingale with characteristics $\left(B^{h}, C, v\right)$. Then, by [JS03, Thm. II.2.34], we have

$$
X=X_{0}+X^{c}+h *\left(\mu^{X}-v\right)+(y-h) * \mu^{X}+B^{h}
$$

### 2.2. Markov processes and related concepts.

2.2.1. Shift semi-group. We consider the shift semi-group

$$
Y_{s} \circ \theta_{t}=Y_{t+s}, \quad t, s \in[0,1] .
$$

2.2.2. Time reversal. We consider the canonical time-reversed process

$$
X_{t}^{*}:=X_{1-t}, \quad t \in[0,1] .
$$

We write $\mathfrak{F}^{*}$ for the associated filtration.
2.2.3. Markov process. A process $\mathrm{P} \in \mathscr{P}(\Omega)$ is Markov provided,

$$
\mathrm{P}\left(X_{t+s} \in B \mid \mathfrak{F}_{t}\right)=\mathrm{P}\left(X_{s} \in B \mid X_{t}\right)
$$

We shall use the following less common alternative characterization. The process $\mathrm{P} \in$ is Markov if and only if

$$
\mathrm{P}\left(C \cap C^{*} \mid X_{t}\right)=\mathrm{P}\left(C \mid X_{t}\right) \mathrm{P}\left(C^{*} \mid X_{t}\right), \quad t \in[0,1], C \in \mathfrak{F}_{t}, C^{*} \in \mathfrak{F}_{t}^{*}
$$

2.2.4. Additive functionals. An adapted process $A$ is an additive functional provided $A_{0}=0$ and

$$
A_{t+s}=A_{s}+A_{t} \circ \theta_{s} .
$$

We identify every additive functional with the random additive set function by letting for $s<t$ :

$$
\begin{aligned}
& A_{[s, t]}:=A_{t-s} \circ \theta_{s}, \\
& A_{[s, t)}:=A_{(t-s)^{-}} \circ \theta_{s}, \\
& A_{(s, t]}:=A_{(t-s)} \circ \theta_{s^{-}}, \\
& A_{(s, t)}:=A_{(t-s)^{-}} \circ \theta_{s^{-}} .
\end{aligned}
$$

2.2.5. Reciprocal measure. We say that P is reciprocal provided

$$
\mathrm{P}\left(X_{s} \in B \mid \mathfrak{F}_{r}, \mathfrak{F}_{t}^{*}\right)=\mathrm{P}\left(X_{s} \in B \mid X_{r}, X_{t}\right), \quad r \leq s \leq t
$$

As for Markov processes, we repeatedly use the alternative characterization: P is reciprocal if and only if

$$
\begin{aligned}
& 0 \leq s \leq t \leq 1, C \in \mathfrak{F}_{s}, C^{*} \in \mathfrak{F}_{t}^{*}, A \in \sigma\left(X_{u}: s \leq u \leq t\right) \Rightarrow \\
& \mathrm{P}\left(C \cap A \cap C^{*} \mid X_{s}, X_{t}\right)=\mathrm{P}\left(C \cap A \mid X_{s}, X_{t}\right) \mathrm{P}\left(C^{*} \cap A \mid X_{s}, X_{t}\right) .
\end{aligned}
$$

Clearly, all Markov processes are also reciprocal. We have that whenever P is reciprocal, then $\mathrm{P}^{x}$ is Markov.
2.3. Feller processes and their generators. We follow [BSW13].
2.3.1. Feller semi-groups. A Feller semi-group is a Markov semi-group ( $\mathbf{T}_{t}$ ) such that

$$
\left\|\mathbf{T}_{t} u-u\right\|_{\infty} \underset{t \rightarrow 0}{ } 0, \quad u \in \mathscr{C}_{0}\left(\mathbb{R}^{d}\right)
$$

2.3.2. Infinitesimal generator. The infinitesimal generator of a Feller semi-group $\left(\mathbf{T}_{t}\right)$ is the unique unbounded operator $\mathbf{A}$ with

$$
\begin{aligned}
& \mathscr{D}(\mathbf{A}):=\left\{u \in \mathscr{C}_{0}\left(\mathbb{R}^{d}\right): \lim _{t \rightarrow 0} \frac{\mathbf{T}_{t} u-u}{t} \text { exists }\right\}, \\
& \mathbf{A} u:=\lim _{t \rightarrow 0} \frac{\mathbf{T}_{t} u-u}{t}
\end{aligned}
$$

The generator $\mathbf{A}$ is densely-defined and closed.
2.3.3. Feller processes and martingales. Let $\mathrm{P} \in \mathscr{P}(\Omega)$ be the law of a Feller process associated with the Feller semi-group ( $\mathrm{T}_{t}$ ), that is

$$
\mathbf{T}_{t} u(x)=\mathbf{E}_{\mathrm{P}}\left[u\left(X_{t}\right) \mid X_{0}=x\right], \quad x \in \mathbb{R}^{d}, t \in[0,1]
$$

Then, for all $f \in \mathscr{D}(\mathbf{A})$, the process

$$
M_{t}:=f\left(X_{t}\right)-f\left(X_{0}\right)-\int_{0}^{t} \mathbf{A} f\left(X_{s}\right) \mathrm{d} s, \quad t \in[0,1]
$$

is a P -martingale.
2.3.4. Lévy triplet of a Feller process. In the rest of the paper, we always assume that $\mathscr{D}(\mathbf{A})$ contains the smooth functions $\mathscr{C}_{c}^{\infty}\left(\mathbb{R}^{d}\right)$. In this case, by [Cou67], there exist:
(i) $\alpha: \mathbb{R}^{d} \rightarrow \mathbb{R}_{+}$;
(ii) $b^{h}: \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}$;
(iii) $c: \mathbb{R}^{d} \rightarrow \mathbb{R}^{d \times d}$ symmetric non-negative;
(iv) $N: \mathbb{R}^{d} \rightarrow \mathcal{M}\left(\mathbb{R}^{d} \backslash\{0\}\right)$ with $\int_{\mathbb{R}^{d}}\left(|y|^{2} \wedge 1\right) N(x, \mathrm{~d} y)<\infty$;
such that

$$
\begin{align*}
\mathbf{A} u(x) & =-\alpha(x) u(x)+b^{h}(x) \cdot \nabla u(x)+\frac{1}{2} \nabla^{2} u(x) \cdot c(x) \\
& +\int_{\mathbb{R}^{d} \backslash\{0\}}(u(x+y)-u(x)-\nabla u(x) \cdot h(y)) N(x, \mathrm{~d} y) . \tag{2.2}
\end{align*}
$$

2.3.5. Feller processes and semimartingales. Let $\mathrm{P} \in \mathscr{P}(\Omega)$ be the law of a Feller process with Lévy triplet $\left(b^{h}, c, N\right)$. Then P is a semimartingale with characteristics

$$
\begin{align*}
& B_{t}^{h}=\int_{0}^{t} b^{h}\left(X_{s}\right) \mathrm{d} s \\
& C_{t}=\int_{0}^{t} c\left(X_{s}\right) \mathrm{d} s  \tag{2.3}\\
& v(\mathrm{~d} s, \mathrm{~d} y)=N\left(X_{s}, \mathrm{~d} y\right) \mathrm{d} s
\end{align*}
$$

### 2.4. Examples of Feller semi-groups.

2.4.1. Heat semi-group. Let R be the law of the Wiener process on $\mathbb{R}^{d}$. Consider the associated Feller semi-group:

$$
\mathbf{T}_{t} u(x):=\int_{\mathbb{R}^{d}} \mathrm{e}^{-|y-x|^{2} / 2 t} \frac{\mathrm{~d} y}{(2 \pi t)^{d / 2}}
$$

Then $\mathbf{A}=-\frac{1}{2} \Delta$. The invariant measure is the Lebesgue measure.
2.4.2. Poisson semi-group. Let $\mathrm{R} \in \mathscr{P}(\Omega)$ be the law of the Poisson process on $\mathbb{R}^{d}$ with intensity $\lambda>0$ and jumping towards $z \in \mathbb{R}^{d}$. The associated Feller semi-group is

$$
\mathbf{T}_{t} u(x):=\sum_{k \in \mathbb{N}} u(x+k z) \frac{(\lambda t)^{j}}{j!} \mathrm{e}^{-t \lambda}
$$

Its generator is $\mathbf{A} u(x)=\lambda(u(x+z)-u(x))$. The invariant measure is the counting measure.
2.4.3. Symmetric stable semi-groups. Let R be the law of the $\alpha$-stable symmetric, with $\alpha \in(0,2)$. The associated Feller semi-group is

$$
\mathbf{T}_{t} u(x):=\int u(x+y) p_{\alpha, t}(\mathrm{~d} y)
$$

where $p_{\alpha, t}$ satisfies

$$
\widehat{p}_{\alpha, t}(\xi)=\mathrm{e}^{-t|\xi|^{\alpha}}
$$

Then the generator is

$$
\mathbf{A} u(x):=k_{\alpha} \int_{\mathbb{R}^{d} \backslash\{0\}}(u(x+y)-u(x)-\nabla u(x) \cdot h(y)) \frac{\mathrm{d} y}{|y|^{\alpha+d}}
$$

2.4.4. Lévy processes. Let R be the law of a Lévy process on $\mathbb{R}^{d}$, that is

$$
\mathbf{T}_{t} u:=u * p_{t}
$$

where $\left(p_{t}\right)$ is a family of infinitely divisible distributions. In this case the Lévy triplet is independent of $x$, and the generator is

$$
\mathbf{A} u(x):=b^{h} \cdot \nabla u(x)+\frac{1}{2} \nabla \cdot(c \nabla u(x))+\int_{\mathbb{R}^{d} \backslash\{0\}}(u(x+y)-u(x)-\nabla u(x) \cdot h(y)) N(\mathrm{~d} y)
$$

2.4.5. Ornstein-Uhlenbeck stable semi-groups. Let R be the law of the $\alpha$-stable symmetric process. The associated Ornstein-Uhlenbeck process $Y$ satisfies, under R

$$
\mathrm{d} Y_{t}=\mathrm{d} X_{t}-Y_{t} \mathrm{~d} t
$$

By solving explicitly this equation, the associated semi-group is

$$
\mathbf{T}_{t} u(x):=\mathbf{E}_{\mathrm{R}}\left[u\left(\mathrm{e}^{t} x+\int_{0}^{t} \mathrm{e}^{s-t} \mathrm{~d} X_{s}\right) \mid X_{0}=x\right]
$$

In this case

$$
\mathbf{A} u(x)=-(-\Delta)^{\alpha / 2} u(x)-x \cdot \nabla u(x)
$$

The unique invariant measure is $\mu_{\alpha}$ such that

$$
\widehat{\mu}_{\alpha}(\xi)=\mathrm{e}^{-\frac{1}{\alpha}|\xi|^{\alpha}}
$$

### 2.5. Relative entropy and Girsanov formula.

2.5.1. Relative entropy. We fix $\mathrm{R} \in \mathscr{P}(\Omega)$ a semimartingale with characteristics $\left(B^{h}, C, \nu\right)$. The relative entropy of with respect to $R$ is the functional

$$
\mathcal{H}(\mathrm{P} \mid \mathrm{R}):= \begin{cases}\mathbf{E}_{\mathrm{P}}\left[\frac{\mathrm{dP}}{\mathrm{dR}}\right], & \text { if } \mathrm{P} \ll \mathrm{R} \\ +\infty, & \text { otherwise }\end{cases}
$$

Elements of $\mathscr{P}(\Omega)$ that are absolutely continuous with respect to to a semimartingale are again a semimartingales [JS03]. In particular, elements with finite entropy with respect to R are semimartingales. [Léo12] characterises their characteristics.
2.5.2. Girsanov theorem under finite entropy. We consider the functions $\theta: \mathbb{R} \ni u \mapsto \mathrm{e}^{u}-u-1$, and its convex conjugate

$$
\theta^{\star}: \mathbb{R} \ni v \mapsto \begin{cases}(v+1) \log (v+1)-v, & \text { if } v>-1 \\ 1, & \text { if } v=-1 \\ +\infty, & \text { otherwise }\end{cases}
$$

Theorem 2.2 ([Léo12, Thms. $1 \& 3]$ ). Then, for all $\mathrm{P} \in \mathscr{P}(\Omega)$ such that $\mathcal{H}(\mathrm{P} \mid \mathrm{R})<\infty$, there exist an adapted process $\beta$, and predictable non-negative generalized process $\ell$ such that

$$
\mathbf{E}_{\mathrm{p}} \int_{0}^{1} \beta_{s} \cdot C(\mathrm{~d} s) \beta_{s}+\mathbf{E}_{\mathrm{p}} \int_{0}^{1} \int_{\mathbb{R}^{d} \backslash\{0\}} \theta^{\star}\left(\left|\ell_{s}(y)-1\right|\right) \nu(\mathrm{d} s \mathrm{~d} y)<\infty .
$$

Moreover, P is a semimartingale with characteristics ( $B^{h}+\tilde{B}, C, \ell \nu$ ), where

$$
\tilde{B}_{t}:=\int_{0}^{t} C(\mathrm{~d} s) \beta_{s}+\int_{0}^{t} \int_{\mathbb{R}^{d} \backslash\{0\}} h(y)\left(\ell_{s}(y)-1\right) v(\mathrm{~d} s \mathrm{~d} y) .
$$

Remark 2.3. [Léo12] states two different theorems, one continuous martingales, and another for pure jump martingales. Decomposing the semimartingale $P$ in its continuous part and pure jump part, it is clear that the results carry out to the mixed case.
2.5.3. Density under finite entropy. Additionally to the Girsanov theorem, [Léo12] obtains an expression of the density in terms of the processes $\beta$ and $\ell$. Consider the log-density process

$$
Z_{t}:=\log \mathrm{E}_{\mathrm{R}}\left[\left.\frac{\mathrm{dP}}{\mathrm{dR}} \right\rvert\, \mathfrak{F}_{t}\right] .
$$

Proposition 2.4 ([Léo12, Thms. $2 \& 4]$ ). On the event $\left\{\frac{\mathrm{dP}}{\mathrm{dR}}>0\right\}$, we have $Z=Z^{c}+Z^{+}+Z^{-}$, where, under P

$$
\begin{aligned}
& Z_{t}^{c}:=\int_{0}^{t} \beta_{s} \cdot\left(\mathrm{~d} X_{s}-\mathrm{d} B_{s}^{h}-C(\mathrm{~d} s) \beta_{s}\right)+\frac{1}{2} \int_{0}^{t} \beta_{s} \cdot C(\mathrm{~d} s) \beta_{s} ; \\
& Z_{t}^{+}:=\int_{[0, t] \times \mathbb{R}_{*}^{d}} \mathbf{1}_{\left\{t \geq \frac{1}{2}\right\}} \log \ell \mathrm{d}\left(\mu^{X}-\ell \nu\right)+\int_{[0, t] \times \mathbb{R}_{*}^{d}} \mathbf{1}_{\left\{\ell \geq \frac{1}{2}\right\}} \theta(\ell-1) \mathrm{d} \nu \\
& Z_{t}^{-}=\int_{[0, t] \times \mathbb{R}_{*}^{d}} \mathbf{1}_{\left\{0 \leq \ell<\frac{1}{2}\right\}} \log \ell \mathrm{d} \mu^{X}+\int_{[0, t] \times \mathbb{R}_{*}^{d}} \mathbf{1}_{\left\{0 \leq t<\frac{1}{2}\right\}}(-\ell+1) \mathrm{d} \nu .
\end{aligned}
$$

Remark 2.5. (a) As noticed in [Léo12], on the event $\left\{\frac{d P}{d R}>0\right\}, \ell>0$ and the sum in the first integral of the definition of $Z^{-}$is finite.
(b) The exact value $1 / 2$ for the cut-off between $Z^{+}$and $Z^{-}$is artificial and could be chosen anywhere in $(0,1)$.
(c) This expression of $Z$ does not provide a decomposition as the sum of $P$-local martingale and absolutely continuous process. Indeed, the stochastic integral $(\log \ell) *\left(\mu^{X}-\ell \nu\right)$ is meaningless, in general. Additional assumptions could guarantee that it makes sense.
2.5.4. Disintegration with respect to the initial condition. So far $\mathrm{P}_{0}$ could be arbitrary, we consider the regular disintegration of P along its marginal $\mathrm{P}_{0}$. In this way, there exists a family $\left(\mathrm{P}^{x}\right)_{x \in \mathbb{R}^{d}}$ such that $\mathrm{P}^{x}$ is supported on $\left\{X_{0}=x\right\}$ and

$$
\mathrm{P}=\int \mathrm{P}^{x} \mathrm{P}_{0}(\mathrm{~d} x) .
$$

The semimartingale property, the Markov property, the special semimartingale property, and the Feller property are stable under this conditioning.
Remark 2.6. In particular, if R is a Markov measure, then all the $\mathrm{R}^{x}$ are also Markov. The data $\left(\Omega, \mathfrak{F},\left(X_{t}\right)_{t \in[0,1]},\left(\mathrm{R}^{x}\right)_{x \in \mathbb{R}^{d}}\right)$ is sometimes what is called a Markov process.
2.5.5. Chain rule for the entropy. By the chain rule for the entropy, [DZ10, Thm. D.13], it turns out that

$$
\begin{equation*}
\mathcal{H}(\mathrm{P} \mid \mathrm{R})=\mathcal{H}\left(\mathrm{P}_{0} \mid \mathrm{R}_{0}\right)+\int \mathcal{H}\left(\mathrm{P}^{x} \mid \mathrm{R}^{x}\right) \mathrm{P}_{0}(\mathrm{~d} x) \tag{2.4}
\end{equation*}
$$

In particular, if $\mathcal{H}(\mathrm{P} \mid \mathrm{R})<\infty$, and that R is a semimartingale, then all the $\mathrm{P}^{x}\left(x \in \mathbb{R}^{d}\right)$ are also semimartingales.

## 3. THE BRENIER-SCHRÖDINGER PROBLEM WITH RESPECT TO A FELLER SEMIMARTINGALE

### 3.1. Formulation of the problem.

3.1.1. Setting. In the rest of the paper, we fix $R \in \mathscr{P}(\Omega)$ a Feller semimartingale with characteristics as in (2.3). Let $\pi \in \mathscr{P}\left(\mathbb{R}^{n} \times \mathbb{R}^{n}\right)$, that we interpret as a coupling between the initial and the final position. Let $\left(\mu_{t}\right)_{t \in[0,1]} \in \mathscr{P}\left(\mathbb{R}^{n}\right)^{[0,1]}$, that we interpret as the incompressibility condition. We study the Brenier-Schrödinger minimisation problem, with respect to the measure $R$

$$
\begin{equation*}
\inf \left\{\mathcal{H}(\mathrm{P} \mid \mathrm{R}): \mathrm{P} \in \mathscr{P}(\Omega), \forall t \in[0,1], P_{t}=\mu_{t}, P_{01}=\pi\right\} \tag{LBS}
\end{equation*}
$$

3.1.2. General shape of the solutions. By [ACLZ20, Thm. 4.7], every minimizer P in (LBS) is a reciprocal measure of the form

$$
\mathrm{P}=\exp \left(\eta\left(X_{0}, X_{1}\right)+A_{1}\right) \mathrm{R}
$$

for some Borel function $\eta: \mathbb{R}^{d} \times \mathbb{R}^{d} \rightarrow \mathbb{R}$ and an additive functional $A$. Conditioning at $\left\{X_{0}=x\right\}$, we find that

$$
\mathrm{P}^{x}=\exp \left(\eta\left(x, X_{1}\right)+A_{1}\right) \mathrm{R}^{x} .
$$

Using the additive property of $A$, we find that

$$
A_{1}=A_{[0,1]}=A_{t^{-}}+A_{[t, 1]} .
$$

It follows that the log-density process looks like

$$
Z_{t}^{x}:=\log \mathbf{E}_{\mathrm{R}^{x}}\left[\left.\frac{\mathrm{dP}^{x}}{\mathrm{dR}^{x}} \right\rvert\, \mathfrak{F}_{t}\right]=A_{t^{-}}+\log \mathbf{E}_{\mathrm{R}^{x}}\left[\exp \left(x, X_{1}\right)+A_{[t, 1]} \mid X_{t}\right]
$$

In view of this formula, let us define

$$
\begin{equation*}
\psi^{x}(t, z):=\log \mathbf{E}_{\mathbb{R}^{x}}\left[\exp \left(\eta\left(x, X_{1}\right)+A_{[t, 1]}\right) \mid X_{t}=z\right], \quad t \in[0,1], z \in \mathbb{R}^{d} \tag{3.1}
\end{equation*}
$$

With this definition, the above expression for $Z^{x}$ reads

$$
Z_{t}^{x}=A_{t^{-}}+\psi^{x}\left(t, X_{t}\right)
$$

3.2. Non-local Hamilton-Jacobi-Bellman equation. Let us reformulate the main result of the paper with our introduced terminology and give the prove.

Theorem 3.1. Let $\mathrm{P} \in \mathscr{P}(\Omega)$ be a regular solution to (LBS). Then, there exists a function $p:[0,1] \times \mathbb{R}^{d}$, such that the potential $\psi^{x}$ defined in (3.1), is a strong solution to

$$
\partial_{t} \psi_{t}^{x}(z)+\mathrm{e}^{-\psi_{t}^{x}(z)}\left(\mathbf{A} \mathrm{e}_{t}^{x}\right)(z)+\frac{1}{2} \nabla \psi_{s}^{x}(z) \cdot c(z) \nabla \psi_{s}^{x}(z)+p_{t}(z)=0
$$

The proof follows from two different representations of $Z^{x}$ as a special martingale.
3.2.1. Decomposition using Itō's formula. For short, write

$$
J_{s}^{x}(y):=\psi_{s}^{x}\left(X_{s^{-}}+y\right)-\psi_{s}^{x}\left(X_{s^{-}}\right) .
$$

Lemma 3.2. Under $\mathrm{P}^{x}$, the process $\psi^{x}(X)$ is a special martingale. Moreover, we have

$$
\begin{aligned}
\psi_{t}^{x}\left(X_{t}\right) & =\psi_{t}^{x}(x)+\int_{0}^{t} \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot \mathrm{d} X_{s}^{c}+J *\left(\mu-\ell^{x} v\right)+\int_{0}^{t} \partial_{s} \psi_{s}^{x}\left(X_{s^{-}}\right) \mathrm{d} s \\
& +\int \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot \mathrm{d} \bar{B}_{s}+\frac{1}{2} \int_{0}^{t} \nabla^{2} \psi_{s}^{x}\left(X^{s^{-}}\right) \cdot \mathrm{d} C_{s}+\int_{0}^{t}\left[J_{s}(y)-\nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot h(y)\right] v(\mathrm{~d} s \mathrm{~d} y)
\end{aligned}
$$

Proof. Since by assumption, $\mathrm{R}^{x}$ is a Feller semimartingale with characteristics $\left(B^{h}, C, v\right)$, by (2.4) and Theorem 2.2, we find that $\mathrm{P}^{x}$ is a semimartingale. Moreover, there exist an adapted process $\beta^{x}$
and a predictable non-negative process $\ell^{x}$ with characteristics ( $\left.B^{h}+\tilde{B}^{x}, C, \ell^{x} \nu\right)$. For short, let us consider the local martingale $M$, the drift $\bar{B}$, and the generalized processes $J$ and $W$ defined by

$$
\begin{aligned}
& M:=X^{c}+h *\left(\mu-\ell^{x} v\right) \\
& \bar{B}:=B^{h}+B^{x} \\
& J_{s}(y):=\psi_{s}^{x}\left(X_{s^{-}}+y\right)-\psi_{s}^{x}\left(X_{s^{-}}\right) \\
& W_{s}(y):=J_{s}(y)-\nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot h(y)
\end{aligned}
$$

Since, by assumption, $\psi^{x}(X)$ is a special semi-martingale, the same argument as in [JS03, Thm. II.2.42] yields that $W * \mu-W *\left(e^{x} v\right)=W *\left(\mu-e^{x} v\right)$ is a local martingale, and

$$
\begin{aligned}
\psi_{t}^{x}\left(X_{t}\right) & =\psi_{t}^{x}(x)+\int_{0}^{t} \partial_{s} \psi_{s}^{x}\left(X_{s^{-}}\right) \mathrm{d} s+\int_{0}^{t} \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot \mathrm{d} M_{s}+\int \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot \mathrm{d} \bar{B}_{s} \\
& +\frac{1}{2} \int_{0}^{t} \nabla^{2} \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot \mathrm{d} C_{s}+W *\left(\mu-\ell^{x} v\right)+W *\left(\ell^{x} v\right)
\end{aligned}
$$

We conclude by observing that

$$
\int_{0}^{t} \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot \mathrm{d}\left(h *\left(\mu-e^{x} v\right)\right)+W *\left(\mu-\ell^{x} v\right)=J *\left(\mu-e^{x} v\right)
$$

### 3.2.2. Decomposition using Girsanov's theorem.

Lemma 3.3. The process $\left(\log \ell^{x}\right) *\left(\mu-\ell^{x} \nu\right)$ is a well-defined local martingale. Moreover, under $\mathrm{P}^{x}$,

$$
\begin{aligned}
Z_{t}^{x} & =\left(\log \ell^{x}\right) *\left(\mu-\ell^{x} v\right)+\int_{0}^{t} \beta_{s}^{x} \cdot \mathrm{~d} X_{s}^{c} \\
& +\int_{0}^{t} \int \theta^{\star}\left(\ell_{s}^{x}(y)-1\right) N\left(X_{s}, \mathrm{~d} y\right) \mathrm{d} s+\frac{1}{2} \int_{0}^{t} \beta_{s}^{x} \cdot c\left(X_{s}\right) \beta_{s}^{x} \mathrm{~d} s
\end{aligned}
$$

where the first line on the right-hand side is a local martingale, and the second line if a predictable process.

Proof. The regular solution assumption ensures that the measure $\mathrm{P}^{x}$ and $\mathrm{R}^{x}$ are equivalent. The part involving $\beta^{x}$ follows directly from Proposition 2.4. In view of Proposition 2.4, it is sufficient to show that

$$
Z_{t}^{+}+Z_{t}^{-}=\left(\log \ell^{x}\right) *\left(\mu-\ell^{x} v\right)+\theta^{\star}\left(\ell^{x}-1\right) * \nu
$$

Recall that

$$
\begin{align*}
Z_{t}^{+}+Z_{t}^{-}= & \int_{[0, t] \times \mathbb{R}_{*}^{d}} \mathbf{1}_{\ell x \geq \frac{1}{2}} \log \left(\ell^{x}\right) \mathrm{d}\left(\mu-\ell^{x} \nu\right)+\int_{[0, t] \times \mathbb{R}_{*}^{d}} \mathbf{1}_{\ell x \geq \frac{1}{2}} \theta^{\star}\left(\ell^{x}-1\right) \mathrm{d} \nu \\
& +\int_{[0, t] \times \mathbb{R}_{*}^{d}} \mathbf{1}_{0 \leq \ell<\frac{1}{2}} \log \left(\ell^{x}\right) \mathrm{d} \mu+\int_{[0, t] \times \mathbb{R}_{*}^{d}} \mathbf{1}_{0 \leq \ell x<\frac{1}{2}}\left(-\ell^{x}+1\right) \mathrm{d} \nu \tag{3.2}
\end{align*}
$$

Since $e^{x} \mathbf{1}_{\ell x \geq 2} \log \left(\ell^{x}\right)$ and $\ell^{x}\left(\mathbf{1}_{1 / 2 \leq \ell<2} \log \left(\ell^{x}\right)\right)^{2}$ are dominated by $\theta^{*}(|\ell-1|)$, which is P $x$-integrable by Theorem 2.2, thus, by [JS03, II.1.27, p. 72], $\mathbf{1}_{\ell x \geq 1 / 2}\left(\log \ell^{x}\right)$ is a local martingale. The second and the fourth terms have $P^{x}$-integrable variations because they are also dominated by $\theta^{*}(|\ell-1|) * \nu$. Since $Z^{x}$ is a special semimartingale, by [JS03, p. I.2.24], the third term also have $\mathrm{P}^{x}$-integrable variations. This implies that the stochastic integral $\mathbf{1}_{0 \leq \ell<1 / 2} \log (\ell) *\left(\mu^{X}-\ell \nu\right)$ is well-defined and a local martingale. Moreover by [JS03, p. II.1.28], we have

$$
\mathbf{1}_{0 \leq \ell^{x}<1 / 2} \log \left(\ell^{x}\right) *\left(\mu-\ell^{x} v\right)=\mathbf{1}_{0 \leq \ell^{x}<1 / 2} \log \left(\ell^{x}\right) * \mu-\mathbf{1}_{0 \leq \ell^{x}<1 / 2} \log \left(\ell^{x}\right) *\left(\ell^{x} v\right)
$$

The proof is thus complete by compensating the third term in (3.2), and by using that the map $a \mapsto a *\left(\mu-\ell^{x} v\right)$ is linear.

### 3.2.3. Conclusion.

Proof of Theorem 3.1. By the uniqueness of the decomposition of a special semimartingale, comparing the expression for $Z^{x}$ in Lemmas 3.2 and 3.3, we obtain by identification of the local martingale parts

$$
\begin{equation*}
\log \ell_{s}^{x}(y)=\psi_{s}^{x}\left(X_{s^{-}}+y\right)-\psi_{s}^{x}\left(X_{s^{-}}\right), \quad \text { and } \quad \beta_{s}^{x}=\nabla \psi_{\cdot}^{x}\left(X_{s^{-}}\right) \tag{3.3}
\end{equation*}
$$

By identification of the predictable part, we obtain

$$
\begin{aligned}
\theta^{\star}\left(\ell^{x}-1\right) * \nu_{s} & +\frac{1}{2} \int_{0}^{t} \beta_{s}^{x} \cdot C(\mathrm{~d} s) \beta_{s}^{x}=A_{s^{-}}+\psi_{0}^{x}(x)+\int_{0}^{t} \partial_{s} \psi_{s}^{x}\left(X_{s^{-}}\right) \mathrm{d} s \\
& +\frac{1}{2} \int_{0}^{t} \nabla^{2} \psi_{s}^{x}\left(X^{s^{-}}\right) \cdot \mathrm{d} C_{s}+\int_{0}^{t} \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot \mathrm{d} B_{s}^{h}+\int_{0}^{t} \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) C(\mathrm{~d} s) \beta_{s}^{x} \\
& +\int_{[0, t]^{\prime} * \mathbb{R}^{d} \backslash\{0\}} \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot h \mathrm{~d}\left(\ell^{x}-1\right) v+\int_{[0, t] \times \mathbb{R}_{*}^{d}}\left[J_{s}-\nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot h\right] \mathrm{d}\left(\ell^{x} v\right) .
\end{aligned}
$$

Reporting (3.3) in the above, terms simplify and we arrive at

$$
\begin{aligned}
& -\int_{[0, t] \times \mathbb{R}_{*}^{d}}\left[\mathrm{e}^{J_{s}}-1-\nabla \psi_{s}^{x}\left(X_{s}\right) \cdot h(y)\right] \mathrm{d} v-\frac{1}{2} \int_{0}^{t} \nabla^{2} \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot \mathrm{d} C_{s}-\frac{1}{2} \int_{0}^{t} \beta_{s}^{x} \cdot C(\mathrm{~d} s) \beta_{s}^{x} \\
& =A_{t^{-}}+\psi_{0}^{x}(x)+\int_{0}^{t} \partial_{s} \psi_{s}^{x}\left(X_{s^{-}}\right) \mathrm{d} s+\int_{0}^{t} \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot \mathrm{d} B_{s}^{h}
\end{aligned}
$$

Recalling that the characteristics have the particular given in (2.3), we see, on the one hand, that $A$ is actually absolutely continuous. Since $A$ is an absolutely continuous additive functional there exist, there exists a pressure $p:=[0,1] \times \mathbb{R}^{d}$ such that

$$
A_{t}=\int_{0}^{t} p_{s}\left(X_{s}\right) \mathrm{d} s
$$

On the other hand, by differentiating with respect to $t$, we obtain

$$
\begin{aligned}
& \partial_{s} \psi_{s}^{x}\left(X_{s^{-}}\right)+p_{s}\left(X_{s}\right)=-\frac{1}{2} \nabla^{2} \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot c\left(X_{s}\right)-\frac{1}{2} \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot c\left(X_{s}\right) \nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \\
& -\mathrm{e}^{-\psi_{t}^{x}\left(X_{t}\right)} \int_{\mathbb{R}_{*}^{d}}\left[\mathrm{e}^{\psi_{t}^{x}\left(X_{t}+y\right)}-\mathrm{e}^{\psi_{t}^{x}\left(X_{t}\right)}-\nabla \mathrm{e}^{\psi_{t}^{x}\left(X_{t}\right)} \cdot h(y)\right] N\left(X_{s}, \mathrm{~d} y\right)+\nabla \psi_{s}^{x}\left(X_{s^{-}}\right) \cdot b^{h}\left(X_{s^{-}}\right)
\end{aligned}
$$

This last equation is true $P^{x}$ almost surely. To conclude from there, we use a continuity argument and the equivalence of $\mathrm{P}^{x}$ and $\mathrm{R}^{x}$.

## 4. EXISTENCE OF SOLUTIONS FOR THE ORNSTEIN-UHLENBECK PROBLEM

Since the functional $\mathcal{H}(\cdot \mid \mathrm{R})$ is convex and lower semi-continuous on $\mathscr{P}(\Omega)$, by the direct method of the calculus of variations, we obtain immediately the following result.

Lemma 4.1. Assume that there exist $\mathrm{P} \in \mathscr{P}(\Omega)$, such that $\mathrm{P}_{01}=\pi, \mathrm{P}_{t}=\mu_{t}$ for all $t \in[0,1]$, and $\mathcal{H}(\mathrm{P} \mid$ $\mathrm{R})<\infty$. Then, there exists a unique minimizer to (LBS).

Thus, in relation with Theorem 3.1, we have to answer two questions:
(A) Can we find a candidate $\mathrm{P} \in \mathscr{P}(\Omega)$ for (LBS) with finite entropy?
(B) Is the unique solution to (LBS) regular?
[ACLZ20, Prop. 6.1] gives a positive answer to (A), when $R$ is the reversible Brownian motion on the torus, $\mu_{t}=$ vol for all $t \in[0,1]$ (incompressible case), and $\pi$ is any coupling satisfying $\mathcal{H}(\pi \mid$ $\left.\mathrm{R}_{01}\right)<\infty$. [GH22] studies the case of the reflected Brownian measure on some quotient spaces, and of a non-reversible Brownian measure in $\mathbb{R}^{n}$. In the latter case, the incompressible condition translates to Gaussian marginal constraints. All the results motioned above, the necessary and sufficient condition for existence is $\mathcal{H}\left(\pi \mid \mathrm{vol}^{2}\right)<+\infty$. In the Ornstein-Uhlenbeck case, we only manage to prove existence when $\pi$ is a Gaussian of a certain form.

In this section, we obtain a positive answer to (A) when R is a reversible Brownian OrnsteinUhlenbeck process on $\mathbb{R}$. In the language of semimartingales, R as characteristics $\left(-X_{t}, \sqrt{2}, 0\right)$, and is started from a Gaussian distribution. In terms of stochastic differential equations, under R , the canonical process $X$ satisfies

$$
\left\{\begin{array}{l}
\mathrm{d} X_{t}=\sqrt{2} \mathrm{~d} W_{t}-X_{t} \mathrm{~d} t \\
X_{0} \sim \gamma:=\mathcal{N}(0,1)
\end{array}\right.
$$

where $W$ is an Brownian motion, under R , independent of $X_{0}$. The measure $\gamma$ is the unique invariant measure of this process, thus we study (LBS) under the natural incompressible condition $\mu_{t}:=\gamma$ for all $t \in[0,1]$. Hence, we consider the minimisation problem Namely, our minimisation problem in this specific case only depends on the parameter $\pi \in \mathscr{P}\left(\mathbb{R}^{n} \times \mathbb{R}^{\mathbb{m}}\right)$, and reads

$$
\begin{equation*}
\inf \left\{\mathcal{H}(\mathrm{P} \mid \mathrm{R}): \mathrm{P} \in \mathscr{P}(\Omega), \mathrm{P}_{t}=\gamma \forall t \in[0,1], \mathrm{P}_{01}=\pi\right\} . \tag{4.1}
\end{equation*}
$$

For $|c| \leq 1$, let us write

$$
\gamma_{c}:=\mathcal{N}\left(0,\left(\begin{array}{ll}
1 & c \\
c & 1
\end{array}\right)\right)
$$

Our results in this case is as follows.
Proposition 4.2. Let $\pi:=\gamma_{c}$ with $4 e^{-1}-3 e^{-1 / 3} \leq c<1$. Then, the problem (4.1) admits a unique solution.

Remark 4.3. Since $\mathrm{P}_{0}=\mathrm{P}_{1}=\gamma$, we necessarily have that the variance under $\pi$ is 1 . Also if $|c|:=1$, then $\pi$ is degenerated and $\mathcal{H}\left(\pi \mid \mathrm{R}_{01}\right)$ is not finite and so the problem would not have any solution. Lastly, for the particular $c:=e^{-1}$, since $\pi=\mathrm{R}_{01}$, the problem admits the trivial solution $P=R$.

As in [ACLZ20; GH22], we create candidate path measures as mixture of R-bridges. In the setting of the Ornstein-Uhlenbeck process, we exploit the following explicit representation for the bridge

$$
\mathrm{R}^{x y}:=\mathrm{R}\left(\cdot \mid X_{0}=x, X_{T}=y\right), \quad x, y \in \mathbb{R}, T \in[0,1] .
$$

Lemma 4.4 ([BK13]). The Ornstein-Uhlenbeck bridge $\mathrm{R}^{x y}$ coincides with the law of the process

$$
U_{t}^{x, y}:=\frac{\sinh (T-t)}{\sinh (T)} x+\frac{\sinh (t)}{\sinh (T)} y+\sqrt{2} \int_{0}^{t} \frac{\sinh (T-t)}{\sinh (T-s)} \mathrm{d} W_{s},
$$

where $W$ is a standard Brownian motion. In particular,

$$
\mathrm{R}_{t}^{x y}=\mathcal{N}\left(\frac{\sinh (T-t)}{\sinh (T)} x+\frac{\sinh (t)}{\sinh (T)} y, 2 \frac{\sinh (T-t) \sinh (t)}{\sinh (T)}\right)
$$

Let $T>0$ and $\sigma \in \mathscr{P}\left(\mathbb{R}^{2}\right)$. We define

$$
\begin{equation*}
\mathrm{Q}:=\int_{\mathbb{R}^{2}} \mathrm{R}^{x y} \sigma(\mathrm{~d} x \mathrm{~d} y) . \tag{4.2}
\end{equation*}
$$

The path measure Q is a mixture of Ornstein-Uhlenbeck bridges.
Remark 4.5. For Brownian bridges, [ACLZ20; GH22] can choose for $\sigma$ a product measure such that the mixture Q satisfies the incompressibility condition, that is $\mathrm{Q}_{t}=$ vol or $\mathrm{Q}_{t}=\mathcal{N}(0,1 / 4)$, for all $t \in[0,1]$. For Ornstein-Uhlenbeck bridges, choosing $\sigma$ as a product cannot yield an invariant process Q. This explains why we need to introduce correlations, and why we are this limited to Gaussian couplings for $\pi$.

Lemma 4.6. Consider the bridge mixture Q as defined in (4.2), with $\sigma:=\gamma_{\rho}$ for some $|\rho|<1$. Then, $\mathrm{Q}_{t}=\gamma$ for all $0 \leq t \leq T$ if and only if $\rho=\mathrm{e}^{-T}$.

Proof. If $\rho=e^{-T}$, then $\sigma=\mathrm{R}_{0 T}$ and $\mathrm{Q}=\mathrm{R}$. Let us show that it is the only possible $\rho$. Let $0 \leq t \leq T$. According to Lemma 4.4 and the definition of $\mathrm{Q}, \mathrm{Q}_{t}$ is the law of

$$
\frac{\sinh (T-t)}{\sinh (T)} X+\frac{\sinh (t)}{\sinh (T)} Z+\sqrt{\frac{\sinh (T-t) \sinh (t)}{\sinh (T)}} W
$$

where $(X, Z) \sim \gamma_{\rho}$ and $W$ is an independent standard Gaussian random variable. In particular, we find for the variance

$$
\operatorname{Var}\left[Q_{t}\right]=\frac{\sinh ^{2}(T-t)}{\sinh ^{2}(T)}+2 \rho \frac{\sinh (T-t) \sinh (t)}{\sinh ^{2}(T)}+\frac{\sinh ^{2}(t)}{\sinh ^{2}(T)}+2 \frac{\sinh (T-t) \sinh (t)}{\sinh (T)}
$$

So, the variance is constant and equals 1 if and only if for all $0<t<T$, we have

$$
2 \rho \sinh (T-t) \sinh (t)=\sinh ^{2}(T)-\sinh ^{2}\left(T_{t}\right)-\sinh ^{2}(t)-2 \sinh (T) \sinh (T-t) \sinh (t)
$$

By direct computations, the right-hand side becomes

$$
2 \sinh (T-t) \sinh (t)(\cosh (T)-\sinh (T))
$$

Thus, the variance is constant and equals to 1 if and only if $\rho=\cosh (T)-\sinh (T)=\mathrm{e}^{-T}$.
Proof of Proposition 4.2. Actually, we need to concatenate several bridges in order to conclude. In this way, we obtain a free parameter for us to optimise. We let $r:=e^{-1 / 3}$, and $s \in \mathbb{R}$ to be chosen later. Let $\sigma \in \mathscr{P}\left(\mathbb{R}^{4}\right)$ be the centred Gaussian law with covariance

$$
C:=\left(\begin{array}{llll}
1 & r & s & c \\
r & 1 & r & s \\
s & r & 1 & r \\
c & s & r & 1
\end{array}\right)
$$

and $\mathrm{Q} \in \mathscr{P}(\Omega)$ defined by

$$
\mathrm{Q}:=\int_{\mathbb{R}^{3}} \mathrm{R}\left(\cdot \mid X_{0}=x, X_{1 / 3}=u, X_{2 / 3}=v, X_{1}=y\right) \sigma(\mathrm{d} x \mathrm{~d} u \mathrm{~d} v \mathrm{~d} y)
$$

The measure $Q$ has finite relative entropy. By the chain rule for the entropy (2.4), we have

$$
\mathcal{H}(\mathrm{Q} \mid \mathrm{R})=\mathcal{H}\left(\pi \mid \mathrm{R}_{01}\right)+\int \mathcal{H}\left(\sigma^{x y} \mid \mathrm{R}_{1 / 3,2 / 3}^{x y}\right) \pi(\mathrm{d} x \mathrm{~d} y)
$$

Since, $\pi$ and $\mathrm{R}_{01}$ on the one hand, and $\sigma^{x y}$ and $\mathrm{R}_{1 / 3,2 / 3}^{x y}$ on the other hand, are non-degenerated Gaussian laws, their relative entropies are finite. Furthermore, $\mathcal{H}\left(\sigma^{x y} \mid \mathrm{R}_{1 / 3,2 / 3}^{x y}\right)$ is a quadratic polynomial in $x$ and $y$. Thus it is integrable with respect to the Gaussian measure $\pi$.

The measure Q satisfies the marginal conditions. By construction we have that $\mathrm{Q}_{01}=\pi$. Let $0<t<1$. Since $R$ is a reciprocal measure, whenever $h \in\{0,1 / 3,2 / 3\}$

$$
\mathrm{Q}_{t}=\int \mathrm{R}_{t}\left(\cdot \mid X_{h}=x, X_{h+1 / 3}=y\right) \gamma_{r}(\mathrm{~d} x \mathrm{~d} y)
$$

Hence, using Lemma 4.6, we have $\mathrm{Q}_{t}=\gamma$.
Handling the parameters. To conclude, let us derive conditions on $s$ and $c$, under which $C$ is a covariance matrix, that is positive definite. Since $C$ is a Toeplitz matrix, its eigenvalues are

$$
\begin{aligned}
& \frac{1}{2}\left(c+r+2 \pm \sqrt{c^{2}-2 c r+5 r^{2}+8 r s+4 s^{2}}\right) \\
& \frac{1}{2}\left(-c-r+2 \pm \sqrt{c^{2}-2 c r+5 r^{2}-8 r s+4 s^{2}}\right)
\end{aligned}
$$

Thus, $C$ is a covariance matrix if and only if

$$
\frac{s^{2}+2 r s+r^{2}-r-1}{r+1}<c<\frac{s^{2}-2 r s+r^{2}+r-1}{r-1}
$$

These two inequalities have solutions if and only if $s \in\left(2 r^{2}-1,1\right)$. Then for each $r\left(4 r^{2}-3\right)<c<1$, there exists $s \in\left(2 r^{2}-1,1\right)$ such that $\Gamma$ is the covariance matrix of a non-degenerated Gaussian measure. This proves the existence of a unique solution.

Remark 4.7. Our candidate measure $Q$ is slightly more involved than the one from [ACLZ20], where the bridge is only conditioned at the time $1 / 2$. In our case, their approach would only prove the existence of solutions for $2 e^{-1}-1 \leq c<1$. Conditioning at times $1 / 3$ and $2 / 3$ gives more flexibility, thanks to the free parameter $s$.

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